DOI: 10.19322/j. cnki. issn. 1006-4710. 2016. 04. 009

曲线坐标系下三维弹性壳体中的微分几何关系

沈晓芹,李昊明

(西安理工大学理学院,陕西西安710054)

摘要:本文建立了三维弹性壳体和其中性面上各点之间的某些微分几何关系表达式,它对形成二 维线性、非线性弹性壳体模型非常重要。具体地,三维弹性体上各点的协变度量张量、逆变度量张 量、度量张量矩阵的行列式以及 Christoffel 符号是由二维中性曲面上的微分几何表达式按壳体厚 度方向的变量渐近展开来表示。

关键词:微分几何;度量张量;Christoffel符号 中图分类号:O186.1 文献标志码:A 文章编号:1006-4710(2016)04-0428-04

Differential geometric relations on the three-dimensional elastic shell in the curvilinear ordinates systems

SHEN Xiaoqin, LI Haoming

(School of Sciences, Xi'an University of Technology, Xi'an 710054, China)

Abstract: The differential geometric relations between 3D elastic shell and the middle surface of shell are provided, which is of importance for forming 2D linear and nonlinear elastic shell models. Concretely, the metric tensor, the determinant of metric matrix field and the Christoffel symbols on the 3D elasticity are expressed by those on the 2D middle surface, which are featured by the asymptotic expressions with respect to the variable in the direction of thickness of the shell.

Key words: differential geometry; metric tensor; Christoffel symbol

在文献[1]和[2]中,三维(3D)区域和二维(2D) 曲面上的微分几何表达式分别定义在曲线坐标上。 在本文中,建立了弹性壳体上各点与壳体中性曲面 之间的微分几何关系式,这种关系式对于从三维方 程推导到 2D 壳体模型非常重要(见文献[3-6])。 笔者借用文献[2]中的符号,下文中,拉丁指标 i,j, $k, \cdots 在集合 {1,2,3} 中取值,而希腊指标 <math>\alpha, \beta, \gamma, \cdots$ 在集合 {1,2}中取值。此外,笔者应用Einstein求和 约定,即重复指标表示求和。矢量 $a,b \in R^3$ 的内积 和外积分别为 $a \cdot b$ 和 $a \times b$ 。

设 $\omega \in R^2$ 上一个有界的连通开子集,其边界 $\gamma = \partial \omega$ 满足 Lipschitz 连续,令 $\gamma = \gamma_0 \cup \gamma_1$,且 $\gamma_0 \cap \gamma_1 = \emptyset_o$ 令 $y = (y_a)$ 表示集合 $\overline{\omega}$ (即 ω 的闭包,见图 1) 中的任意点,且 $\partial_a := \partial/\partial y_a$ 。

设单射 $\boldsymbol{\theta} \in C^{\mathfrak{s}}(\bar{\omega}; \mathbb{R}^{\mathfrak{s}}), \mathbb{M}$: $\boldsymbol{a}_{\mathfrak{a}}(y) := \partial_{\mathfrak{a}} \boldsymbol{\theta}(y)$ (1)



图 1 壳体 $\hat{\Omega}^{\epsilon}$ 及其中性面 $S^{[2]}$ Fig. 1 Shell $\hat{\Omega}^{\epsilon}$ with middle surface $S^{[2]}$

在点 $y \in \omega$ 处是线性无关的。因此这两个矢量在曲面 $S := \theta(\omega)$ 上的点 $\theta(y)$ 处张成其切平面,且单位 矢量 $a_s(y)$ 在点 $\theta(y)$ 处垂直于 S_s

收稿日期: 2016-01-08

基金项目:国家自然科学基金资助项目(11571275);陕西省工业科技攻关资助项目(2015GY021);陕西省教育厅基金资助项目(2015CX009)

作者简介:沈晓芹,女,博士,副教授,研究方向为微分方程数值解及其应用。E-mail: xqshen@xaut.edu.cn

$$\boldsymbol{a}_{3}(y) := \frac{\boldsymbol{a}_{1}(y) \times \boldsymbol{a}_{2}(y)}{|\boldsymbol{a}_{1}(y) \times \boldsymbol{a}_{2}(y)|}$$

这些矢量 $a_i(y)$ 在点 $\theta(y)$ 处构成了协变基矢量,而矢量 $a^i(y)$ 定义为:

$$a^i(y) \cdot a_i(y) = \delta^i$$

在点 $\theta(y)$ 处构成逆变基矢量,其中 δ_j 是 Kronecker 符号(注意 $a^3(y) = a_3(y)$ 且矢量 $a^a(y)$ 在 S 的切平 面上)。

S上的度量张量的协变分支 a_{φ} 和逆变分支 a^{φ} 、S 上的 Christoffel 符号 $\Gamma_{\varphi,\sigma}$ 和 Γ_{φ}^{c} 、S 上的曲率张量的协 变分支 b_{φ} 和混合分支 b_{α}^{s} 、S 上第三基本形的协变分支 定义如下(显然依赖于变量 $y \in \overline{\omega}$,因此略去):

$$\begin{cases}
 a_{\alpha\beta} := a_{\alpha} \cdot a_{\beta} \\
 a^{\alpha\beta} := a^{\alpha} \cdot a^{\beta} \\
 (a^{\alpha\beta}) = (a_{\alpha\beta})^{-1} \\
 a^{\alpha} = a^{\alpha\beta}a_{\beta}
\end{cases}$$
(2)

$$\begin{cases} \stackrel{\circ}{\Gamma}_{\alpha\beta,\sigma} := a_{\sigma} \cdot \partial_{\alpha} a_{\beta} \\ \stackrel{\circ}{\Gamma}_{\sigma} := a^{\sigma} \cdot \partial_{\alpha} a \end{cases}$$
(3)

$$\begin{cases} b_{\alpha\beta} := a_3 \cdot \partial_{\alpha} a_{\beta} \\ b_{\alpha}^{\beta} := a^{\beta\sigma} \cdot b_{\sigma\alpha} \\ c_{\alpha} := \partial a_2 \cdot \partial_{\alpha} a_{\beta} \end{cases}$$
(4)

其中(*a_{qq}*)是对称正定矩阵,(*b_{qq}*)和(*c_{qq}*)是对称矩阵。度量张量、曲率张量和第三基本形的行列式定义为:

$$\begin{cases} a := \det(a_{\alpha\beta}) \\ b := \det(b_{\alpha\beta}) \\ c := \det(c_{\alpha\beta}) \end{cases}$$

假设有一壳体 $\hat{\Omega}^{\epsilon}$ (见图 1)和中性面 $S = \boldsymbol{\theta}(\bar{\omega})$, 其厚度 2 ϵ >0 任意小。因此,对每一个 ϵ >0,壳体的 参考构型为 $\hat{\Omega}^{\epsilon} = \boldsymbol{\Theta}(\overline{\Omega}^{\epsilon})$,其中 $\overline{\Omega}^{\epsilon} = \bar{\omega} \times [\cdot \epsilon, \epsilon]$,即:

$$\boldsymbol{\Theta}(y,\boldsymbol{\xi}) = \boldsymbol{\theta}(y) + \boldsymbol{\xi} \boldsymbol{a}_{3}(y)$$

売 体 $\Theta(\overline{\Omega}^{\epsilon})$ 的 上 表 面 和 下 表 面 分 别 为 $\Gamma_{\iota} = \Theta(\omega \times \{+\epsilon\}), \Gamma_{b} = \Theta(\omega \times \{-\epsilon\})$ 。 売 体 側 面 为 $\Gamma_{\iota} = \Gamma_{0} \cup \Gamma_{1},$ 其中^[7]:

 $\begin{cases} \Gamma_0 = \boldsymbol{\theta}(\boldsymbol{\gamma}_0) \times (-\boldsymbol{\varepsilon}, +\boldsymbol{\varepsilon}) \\ \Gamma_1 = \boldsymbol{\theta}(\boldsymbol{\gamma}_1) \times (-\boldsymbol{\varepsilon}, +\boldsymbol{\varepsilon}) \end{cases}$

令 $x = (x_i)$ 表示集合 $\overline{\Omega}^{\epsilon}$ 上的任一点,其中 $x_a = y_a, x_3 = \xi$ 。映射 $\Theta: \overline{\Omega}^{\epsilon} \to R^3$ 是单射且三个矢量 $g_i(x) := \partial_i \Theta(x)$ 在点 $x \in \overline{\Omega}^{\epsilon}$ 处是线性无关的。 矢量 $g^i(y)$ 定义为:

 $\boldsymbol{g}^{i}(x) \boldsymbol{\cdot} \boldsymbol{g}_{j}(x) = \delta^{i}_{j}$

这些矢量在点 $\Theta(x)$ 处构成了逆变基矢量。

 $\Theta(\overline{\Omega}^{e})$ 上的度量张量的协变分支 g_{ij} 和逆变分

支 g^{ij} 、 $\Theta(\overline{\Omega}^{\epsilon})$ 上的 Christoffel 符号 $\Gamma_{ij,k}$ 和 Γ_{ij}^{k} ,被定 义如下(显然依赖于变量 $x \in \overline{\Omega}^{\epsilon}$,因此略去):

$$\begin{cases} g_{ij} := g_i \cdot g_j \\ g^{ij} := g^i \cdot g^j \\ \begin{cases} \Gamma_{ij,k} := g_k \cdot \partial_i g_j \\ \Gamma_{ij}^k := g^{kl} \Gamma_{ij,l} \end{cases}$$
(5)

度量张量的行列式为:

$$g := \det(g_{ij})$$

在第二部分中,3D区域上的度量张量、度量矩 阵的行列式和 Christoffel 符号可由 2D 曲面上表达 式按壳体厚度方向的变量渐近展开来表示。

1 主要结论

定理1 设壳体 $\Theta(y,\xi)$ 和中性面 $S = \theta(\bar{\omega})$ 上的度量张量分别为 g_{ij} 和 a_{qg} , b_{qg} 和 c_{qg} 是 S 上第二和 第三基本形, 那么下面的微分几何关系式成立。

$$\begin{cases} g_{\alpha\beta} = a_{\alpha\beta} - 2\xi b_{\alpha\beta} + \xi^2 c_{\alpha\beta} \\ g_{\alpha3} = g_{3\alpha} = 0 \\ g_{33} = 1 \end{cases}$$

$$\nexists \oplus \alpha, \beta = 1, 2, \xi \in [-\varepsilon, \varepsilon]_{\circ}$$

$$\blacksquare \blacksquare \blacksquare :$$

$$g_{\alpha\beta} = g_{\alpha} \cdot g_{\beta} = \partial_{\alpha} \oslash \cdot \partial_{\beta} \oslash = \\ \partial_{\alpha} (\varTheta + \xi a_{3}) \cdot \partial_{\beta} (\varTheta + \xi a_{3}) = \\ \partial_{\alpha} \varTheta \cdot \partial_{\beta} \varTheta + \partial_{\alpha} \varTheta \cdot \partial_{\beta} (\xi a_{3}) + \\ \partial_{\alpha} (\xi a_{3}) \cdot \partial_{\beta} \varTheta + \partial_{\alpha} (\xi a_{3}) \cdot \partial_{\beta} (\xi a_{3}) = \\ a_{\alpha} \cdot a_{\beta} + \xi a_{\alpha} \cdot \partial_{\beta} a_{3} + \\ \xi \partial_{\alpha} a_{3} \cdot a_{\beta} + \xi^{2} \partial_{\alpha} a_{3} \cdot \partial_{\beta} a_{3} \end{cases}$$
(6)

将式(1)~(4)代入(6),基于 b_a的对称性,得到:

$$g_{\alpha\beta} = a_{\alpha\beta} - 2\xi b_{\alpha\beta} + \xi^{2} c_{\alpha\beta}$$

$$g_{3a} = g_{3} \cdot g_{a} = \partial_{3} \Theta \cdot \partial_{a} \Theta =$$

$$\partial_{3} (\theta + \xi a_{3}) \cdot \partial_{a} (\theta + \xi a_{3}) =$$

$$a_{3} \cdot \partial_{a} (\theta + \xi a_{3}) =$$

$$a_{3} \cdot \partial_{a} \theta_{3} + a_{3} \cdot \partial_{a} (\xi a_{3}) =$$

$$a_{3} \cdot a_{a} + \xi a_{3} \cdot \partial_{a} a_{3}$$

$$f H = a_{3} \phi h = 0$$
(7)

根据 a_3 的定义,有

$$\begin{cases} \boldsymbol{a}_3 \cdot \boldsymbol{a}_a = 0 \\ \boldsymbol{a}_3 \cdot \boldsymbol{a}_3 = 1 \end{cases}$$
(8)

则:

因

$$\partial_{a} (\boldsymbol{a}_{3} \cdot \boldsymbol{a}_{3}) = 2\boldsymbol{a}_{3} \cdot \partial_{a}\boldsymbol{a}_{3} = 0$$

$$\mathbb{L}:$$

$$\boldsymbol{a}_{3} \cdot \partial_{a}\boldsymbol{a}_{3} = 0 \tag{9}$$

类似地:

$$g_{a3} = 0$$

$$g_{33} = g_3 \cdot g_3 = \partial_3 \Theta \cdot \partial_3 \Theta =$$

$$\partial_3 (\theta + \xi a_3) \cdot \partial_3 (\theta + \xi a_3) =$$

$$\partial_3 \theta \cdot \partial_3 \theta + \partial_3 \theta \cdot \partial_3 (\xi a_3) +$$

$$\partial_3 (\xi a_3) \cdot \partial_3 \theta + \partial_3 (\xi a_3) \cdot \partial_3 (\xi a_3) =$$

$$0 + 0 + 0 + a_3 \cdot a_3 = 1$$

证毕。

 \cap

由于 $(g_{ij}) = (g^{ij})^{-1}$,逆变分支 g^{ij} 的表达式如 定理 2 所示。

定理 2 在定理 1 的假设下,令 $g^{ij} \neq \Theta(y,\xi)$ 上度量张量的逆变分支,则下述关系式成立:

$$\begin{cases} g^{11} = g^{-1} (a_{22} - 2\xi b_{22} + \xi^2 c_{22}) \\ g^{12} = g^{21} = -g^{-1} (a_{12} - 2\xi b_{12} + \xi^2 c_{12}) \\ g^{22} = g^{-1} (a_{11} - 2\xi b_{11} + \xi^2 c_{11}) \\ g^{a3} = g^{3a} = 0 \\ g^{33} = 0 \\ g^{33} = 0 \\ \ddagger \varphi + \vdots \\ g = \det(g_{ij}) = \end{cases}$$
(10)

$$(a_{11} - 2\xi b_{11} + \xi^2 c_{11}) (a_{22} - 2\xi b_{22} + \xi^2 c_{22}) - (a_{12} - 2\xi b_{12} + \xi^2 c_{12})^2$$

证明:

$$(g_{ij})^{-1} = \begin{pmatrix} (g_{q\beta}) & 0 \\ 0 & g_{33} \end{pmatrix}^{-1} = \begin{pmatrix} (g_{q\beta})^{-1} & 0 \\ 0 & (g_{33})^{-1} \end{pmatrix}$$

其中,

$$(g_{aeta})^{-1} = egin{pmatrix} g_{11} & g_{12} \ g_{21} & g_{22} \end{pmatrix}^{-1} = g^{-1} egin{pmatrix} g_{22} & -g_{12} \ -g_{21} & g_{11} \end{pmatrix}$$

由于 $(g^{ij}) = (g_{ij})^{-1}$,可以很容易地导出式(10)。 证毕。

定理3 在定理1的假设下,令 $\Gamma_{ij,k}$, $\stackrel{*}{\Gamma}_{\varphi,\gamma}$ 分别 是 $\Theta(y,\xi)$ 和 $\theta(\bar{\omega})$ 上的Christoffel符号,则式(11) 成立。

$$\begin{cases} \Gamma_{\alpha\beta,\sigma} = \tilde{\Gamma}_{\alpha\beta,\sigma} + \xi a_{\sigma} \cdot \partial_{\alpha\beta} a_{3} + \\ \xi \partial_{\sigma} a_{3} \cdot \partial_{\alpha} a_{\beta} + \xi^{2} \partial_{\sigma} a_{3} \cdot \partial_{\alpha\beta} a_{3} \\ \Gamma_{\alpha\beta,3} = b_{\alpha\beta} - c_{\alpha\beta} \\ \Gamma_{\alpha3,\sigma} = \Gamma_{3\alpha,\sigma} = -b_{\alpha\sigma} + \xi c_{\alpha\sigma} \\ \Gamma_{33,\alpha} = \Gamma_{\alpha3,3} = \Gamma_{3\alpha,3} = \Gamma_{33,3} = 0 \\ \text{It } \oplus \alpha, \beta, \sigma = 1, 2_{\circ} \\ \text{it II:} \\ \Gamma_{\alpha\beta,\sigma} = g_{\sigma} \cdot \partial_{\alpha} g_{\beta} = \partial_{\sigma} \Theta \cdot \partial_{\alpha} (\partial_{\beta} \Theta) = \end{cases}$$

$$\end{cases}$$

$$\begin{aligned} & \mathcal{A}_{\varphi,\sigma} = \mathbf{g}_{\sigma} \cdot \partial_{\alpha} \mathbf{g}_{\beta} = \partial_{\sigma} \mathbf{O} \cdot \partial_{\alpha} \left(\partial_{\beta} \mathbf{O} \right) = \\ & \partial_{\sigma} \left(\mathbf{\theta} + \boldsymbol{\xi} \mathbf{a}_{\beta} \right) \cdot \partial_{\varphi\beta} \left(\mathbf{\theta} + \boldsymbol{\xi} \mathbf{a}_{\beta} \right) = \\ & \partial_{\sigma} \mathbf{0} \cdot \partial_{\varphi\beta} \mathbf{\theta} + \partial_{\sigma} \mathbf{0} \cdot \partial_{\varphi\beta} \left(\boldsymbol{\xi} \mathbf{a}_{\beta} \right) + \\ & \partial_{\sigma} \left(\boldsymbol{\xi} \mathbf{a}_{\beta} \right) \cdot \partial_{\varphi\beta} \mathbf{\theta} + \partial_{\sigma} \left(\boldsymbol{\xi} \mathbf{a}_{\beta} \right) \cdot \partial_{\varphi\beta} \left(\boldsymbol{\xi} \mathbf{a}_{\beta} \right) = \\ & \overset{*}{\Gamma}_{\varphi,\sigma} + \boldsymbol{\xi} \mathbf{a}_{\sigma} \cdot \partial_{\varphi\beta} \mathbf{a}_{\beta} + \end{aligned}$$

$$\begin{aligned} \boldsymbol{\xi}\partial_{\sigma}\boldsymbol{a}_{3} \cdot \partial_{a}\boldsymbol{a}_{\beta} + \boldsymbol{\xi}^{2}\partial_{\sigma}\boldsymbol{a}_{3} \cdot \partial_{q\beta}\boldsymbol{a}_{3} & (12) \\ \text{ lef } \boldsymbol{f} \cdot \boldsymbol{a}_{3} \cdot \boldsymbol{a}_{3} = 1, \boldsymbol{f} \cdot \boldsymbol{f} \\ \partial_{q\beta}\left(\boldsymbol{a}_{3} \cdot \boldsymbol{a}_{3}\right) = \partial_{a}\left(\partial_{\beta}\left(\boldsymbol{a}_{3} \cdot \boldsymbol{a}_{3}\right)\right) = \\ \partial_{a}\left(2\boldsymbol{a}_{3} \cdot \partial_{\beta}\boldsymbol{a}_{3}\right) = \\ 2\partial_{a}\boldsymbol{a}_{3} \cdot \partial_{\beta}\boldsymbol{a}_{3} + 2\boldsymbol{a}_{3} \cdot \partial_{q\beta}\boldsymbol{a}_{3} = 0 \end{aligned}$$

则:

$$a_{3} \cdot \partial_{a\beta}a_{3} = -\partial_{a}a_{3} \cdot \partial_{\beta}a_{3}$$
(13)
将式(13)代人(12),得:
$$\Gamma_{a\beta,3} = b_{a\beta} - c_{a\beta}$$

$$\Gamma_{a3,\sigma} = g_{\sigma} \cdot \partial_{a}g_{3} = \partial_{\sigma}\Theta \cdot \partial_{a}(\partial_{3}\Theta) =$$

$$\partial_{\sigma}(\theta + \xi a_{3}) \cdot \partial_{a}(\partial_{3}(\theta + \xi a_{3})) =$$

$$\partial_{\sigma}(\theta + \xi a_{3}) \cdot \partial_{a}a_{3} =$$

$$\partial_{\sigma}\theta \cdot \partial_{a}a_{3} + \xi \partial_{\sigma}a_{3} \cdot \partial_{a}a_{3} =$$

$$-b_{a\sigma} + \xi c_{a\sigma}$$

类似地,

$$\Gamma_{3a,\sigma} = -b_{a\sigma} + \pounds c_{a\sigma}$$

$$\Gamma_{33,a} = g_a \cdot \partial_3 g_3 = \partial_a \Theta \cdot \partial_3 (\partial_3 \Theta) =$$

$$\partial_a (\theta + \pounds a_3) \cdot \partial_3 (\partial_3 (\theta + \pounds a_3)) =$$

$$\partial_a (\theta + \pounds a_3) \cdot \partial_3 a_3 = 0$$

$$\Gamma_{a3,3} = g_3 \cdot \partial_a g_3 = \partial_3 \Theta \cdot \partial_a (\partial_3 \Theta) =$$

$$\partial_3 (\theta + \pounds a_3) \cdot \partial_a (\partial_3 (\theta + \pounds a_3)) =$$

$$a_3 \cdot \partial_a a_3 = 0$$

$$\Gamma_{33,3} = g_3 \cdot \partial_3 g_3 = \partial_3 \Theta \cdot \partial_3 (\partial_3 \Theta) =$$

$$\partial_3 (\theta + \pounds a_3) \cdot \partial_3 (\partial_3 (\theta + \pounds a_3)) =$$

$$a_3 \cdot \partial_3 a_3 = 0$$
IE 毕 。

因此, Christoffel 符号 $\Gamma_{ij}^{t} \pi \Gamma_{aj}^{r_{o}}$ 有类似的关系。

定理 4 在定理 1 的假设下,令 $\Gamma_{ij}^{\xi} \neq \Theta(y,\xi)$ 上的 Christoffel 符号,则式(14)成立。

$$\begin{cases} \Gamma_{q\beta}^{\tau} = g^{\sigma\tau} (\mathring{\Gamma}_{q\beta,\tau} + \xi a_{\tau} \cdot \partial_{q\beta} a_{3} + \xi^{2} \partial_{\tau} a_{3} \cdot \partial_{q\beta} a_{3}) \\ \Gamma_{q\beta}^{3} = b_{q\beta} - c_{q\beta} \\ \Gamma_{a3}^{\tau} = \Gamma_{3a}^{\tau} = g^{\sigma\tau} (-b_{a\sigma} + \xi c_{a\tau}) \\ \Gamma_{33}^{\tau} = \Gamma_{a3}^{\theta} = \Gamma_{3a}^{\theta} = \Gamma_{33}^{\theta} = 0 \end{cases}$$
(14)
$$\nexists \ \ \mu \ \alpha, \beta, \sigma = 1, 2,$$
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(15)(15)
$$\square \ \mu \ \alpha, \beta,$$

$$\begin{split} \Gamma_{33}^{a} &= g^{al} \Gamma_{33,l} = \\ g^{a\tau} \Gamma_{33,\tau} + g^{a3} \Gamma_{33,3} = 0 \\ \Gamma_{a3}^{3} &= \Gamma_{3a}^{3} = g^{3l} \Gamma_{a3,l} = \\ g^{3\tau} \Gamma_{a3,\tau} + g^{33} \Gamma_{a3,3} = \Gamma_{a3,3} = 0 \\ \Gamma_{33}^{3} &= g^{3l} \Gamma_{33,l} = \\ g^{3\tau} \Gamma_{33,\tau} + g^{33} \Gamma_{33,3} = \Gamma_{33,3} = 0 \end{split}$$

因此,式(14)能够容易地从定理2和定理3而 得到。

证毕。

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